

Molecular Spectroscopy in Chemistry

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Schedule: every Friday,

16:15-18:00, room CH B331

Caldara Tiziano Agostino
Clerc Hugo André
Ding Zixuan
Fischer Thomas
Gaedecke Kelian Dietmar Marc
Hawkins Jasper Raymond
Huber Jan
Hügli Loïc
Jaccard Gladys Gwendoline Geneviève
Mikov Aleksandar Mihaylov
Nasturzio Pietro
Venancio Enzo Franck
Xie Qishan
Zupan Léa

*13 lectures;
one session for solving problems;
oral exam.*

Online course PDF presentations in MOODLE ([CH-456](#));

Course structure

◆ Revision of basics

Light as wave and particle
Quantum mechanics;
Spectroscopic transitions

◆ Types of spectroscopy

Absorption
Fluorescence
Raman

◆ Analytical Spectroscopy in Chemistry

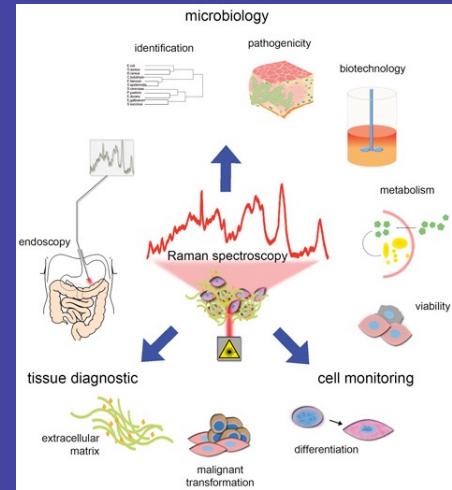
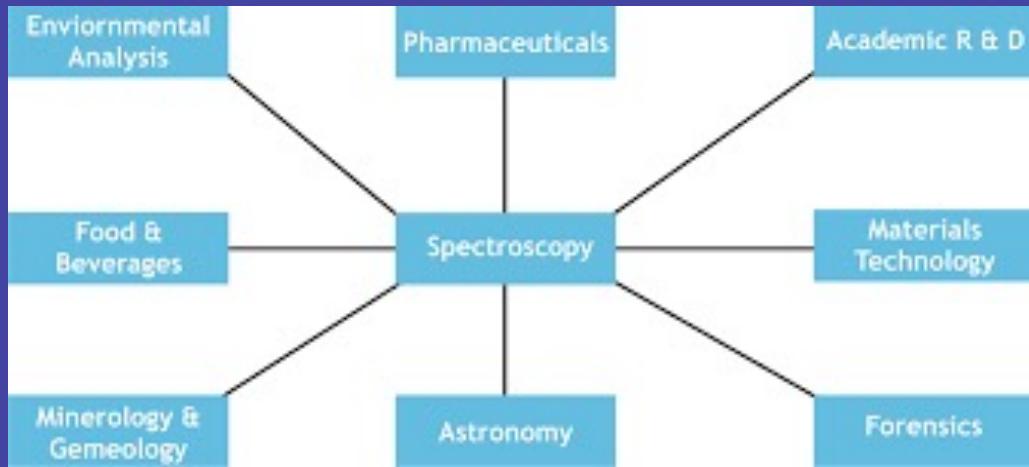
Laser
FTIR
Raman microscopy
Photofragmentation

◆ Selective chemistry and Laser isotope separation

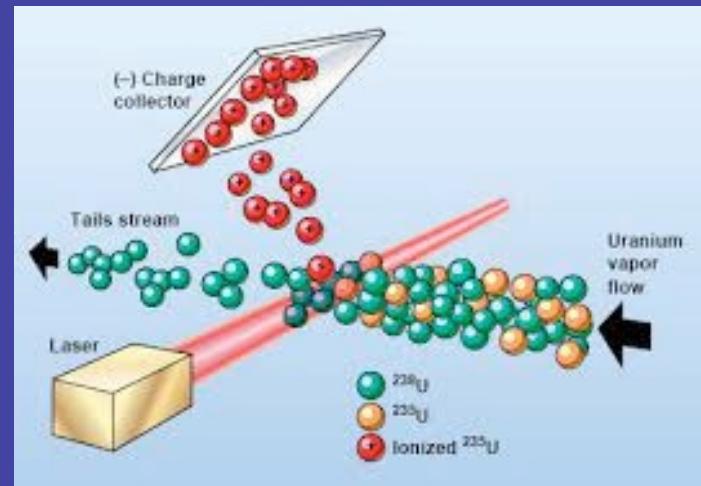
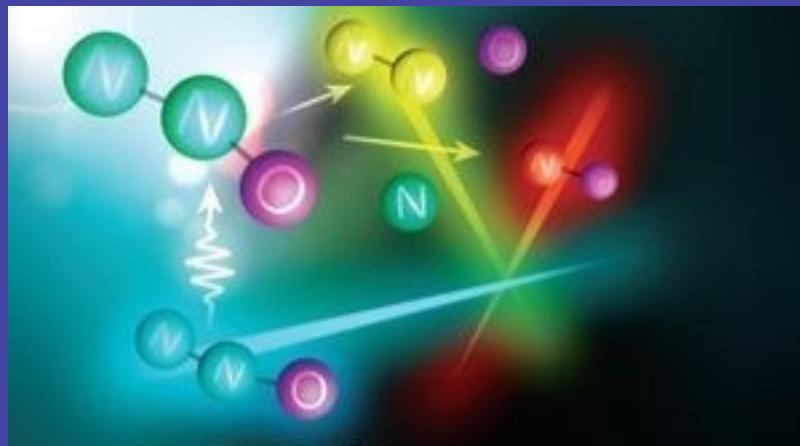
◆ Spectroscopy in Astronomy (Solar and exoplanets)

Spectroscopy in chemistry.

- *Analytical:* what, how much, where ?



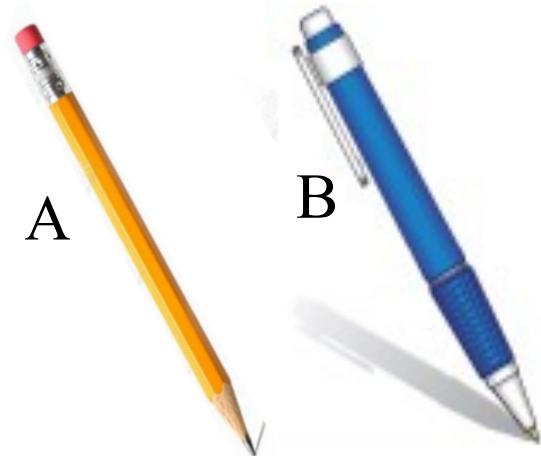
- *Action:* light-driven selective chemistry



Concept of “library”

Which image is for a pen?

- *You have sets of visual properties in your memory for objects;*
- *You brain compares the images with these sets to identify them.*



*Similar for molecules: one needs a set of physical observables to **identify** a molecule: e.g., mass (MS), retention time (LC), chemical shifts (NMR), ...*

For each technique these properties are to be measured (exceptionally computed) in advance for known standards and stored.

*This is called “**Library approach**”*

Spectroscopy identifies a molecule by its spectra: a result of interaction of light with molecules as a function of light wavelength.

Quantification

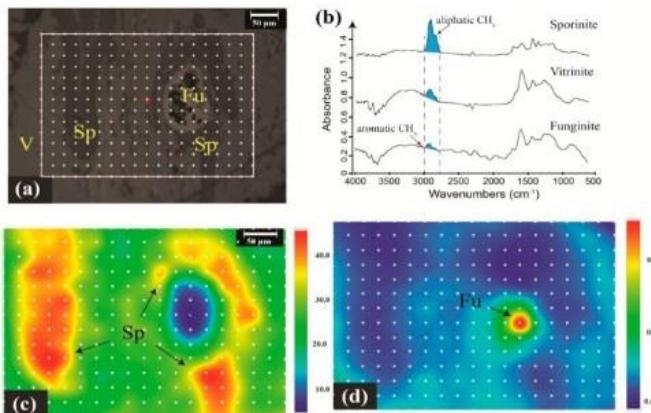
Which box is larger?

- Need a **standard** to compare:
e.g., attenuation of light by a sample relative to a standard of known concentration.

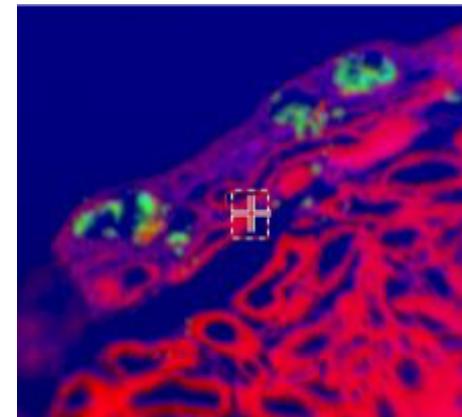


Location

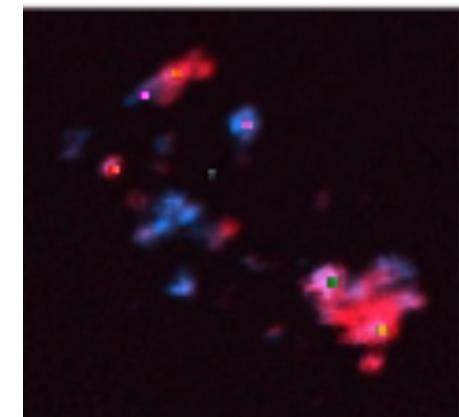
Light is to be directed to or light action is to be detected from a desired volume/surface only



Geological science



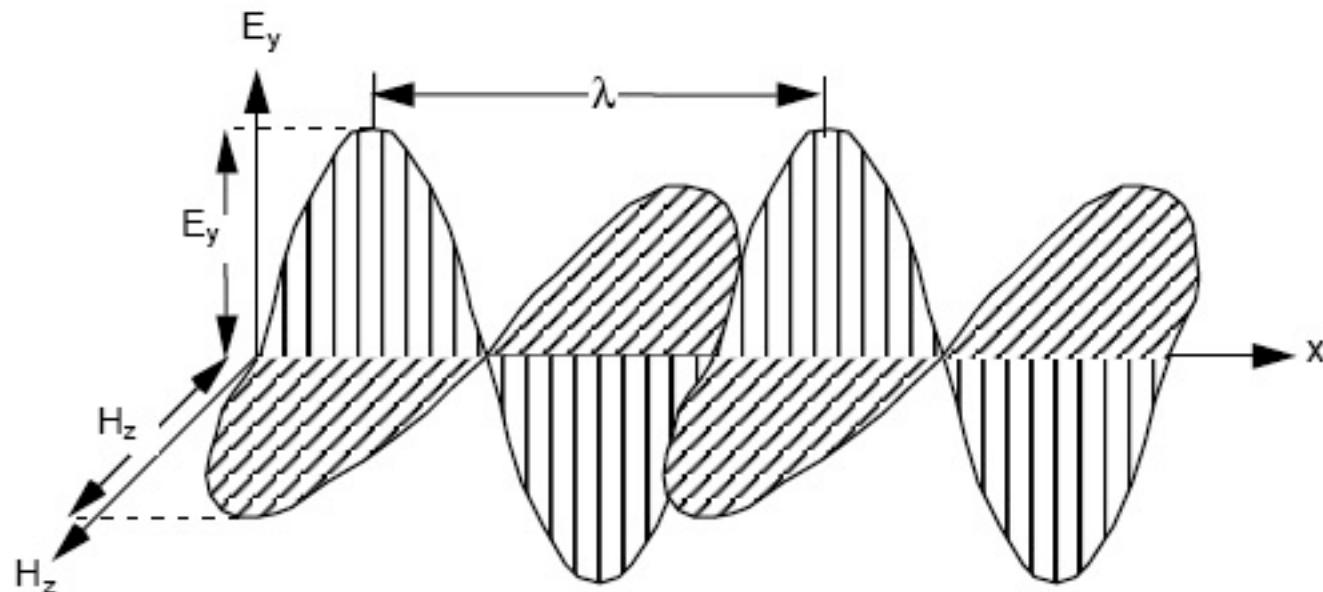
Cancer cells in tissue



Industrial defects

Review of fundamentals

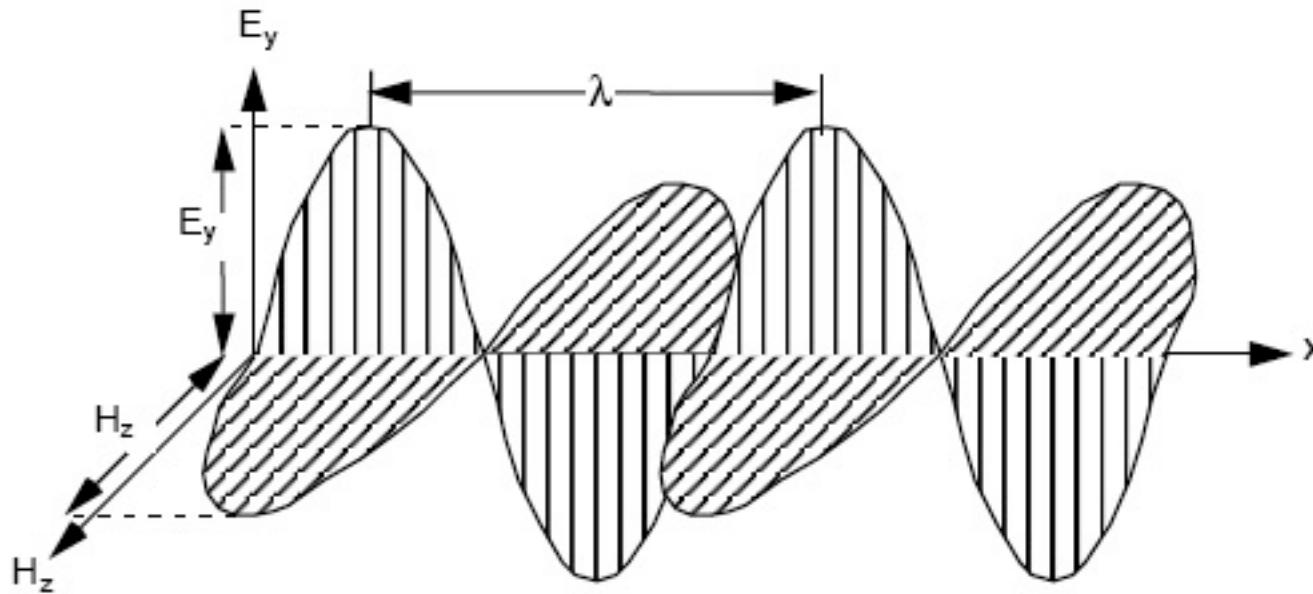
Properties of Electromagnetic Radiation



$$c = \lambda f \cong 3 \cdot 10^8 \text{ m/s}$$

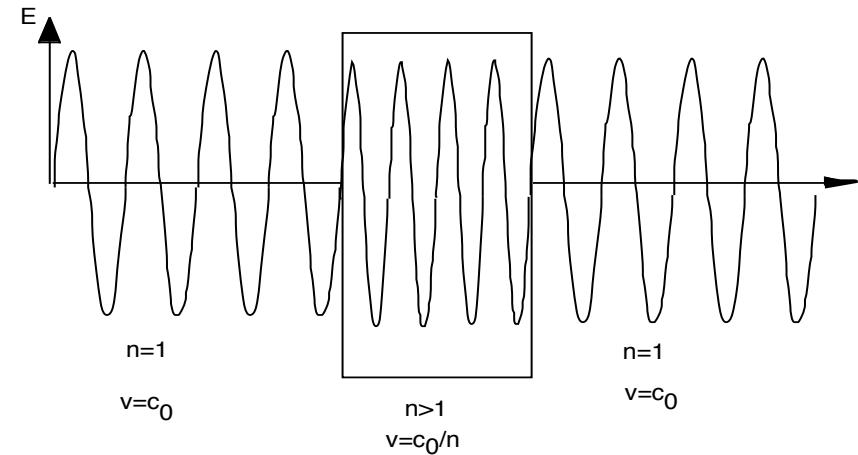
$$\lambda = \frac{c}{f}$$

Properties of Electromagnetic Radiation



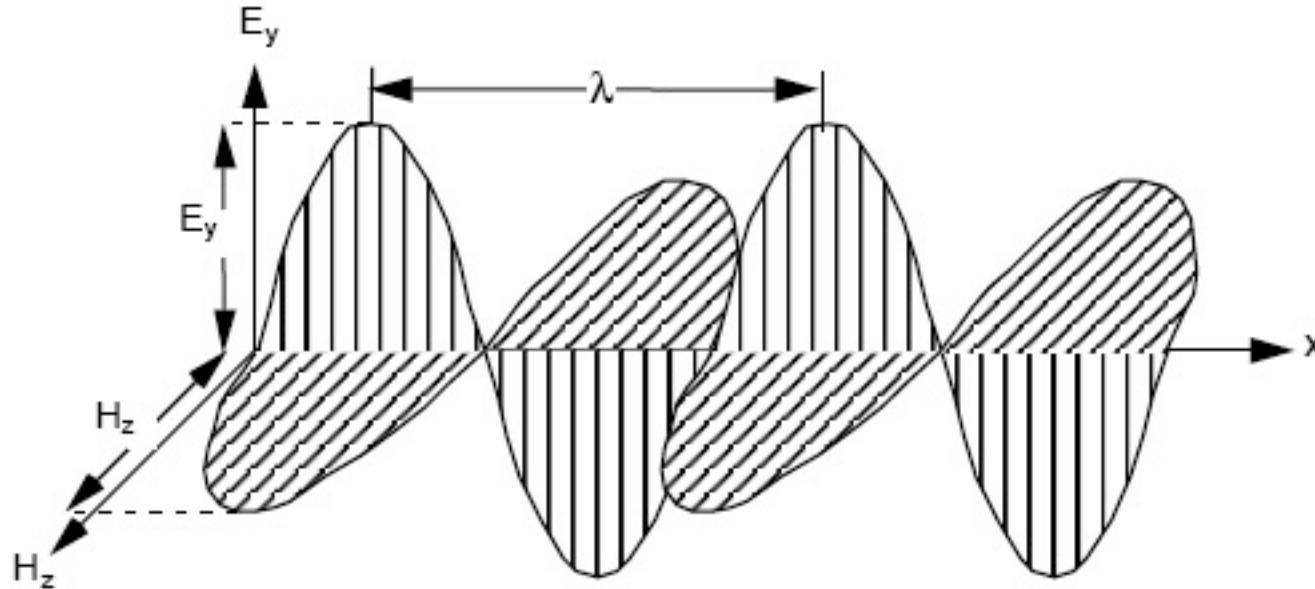
Refractive index of media $n = \frac{c}{V}$

$$v = \frac{c}{n} = \lambda_n \cdot f \quad \lambda_n = \frac{c}{n(f) \cdot f}$$



EMR propagates slower in any media, but with the same frequency

Properties of Electromagnetic Radiation



$$E(x,t) = E_{y0} \sin(2\pi ft + kx + \phi)$$

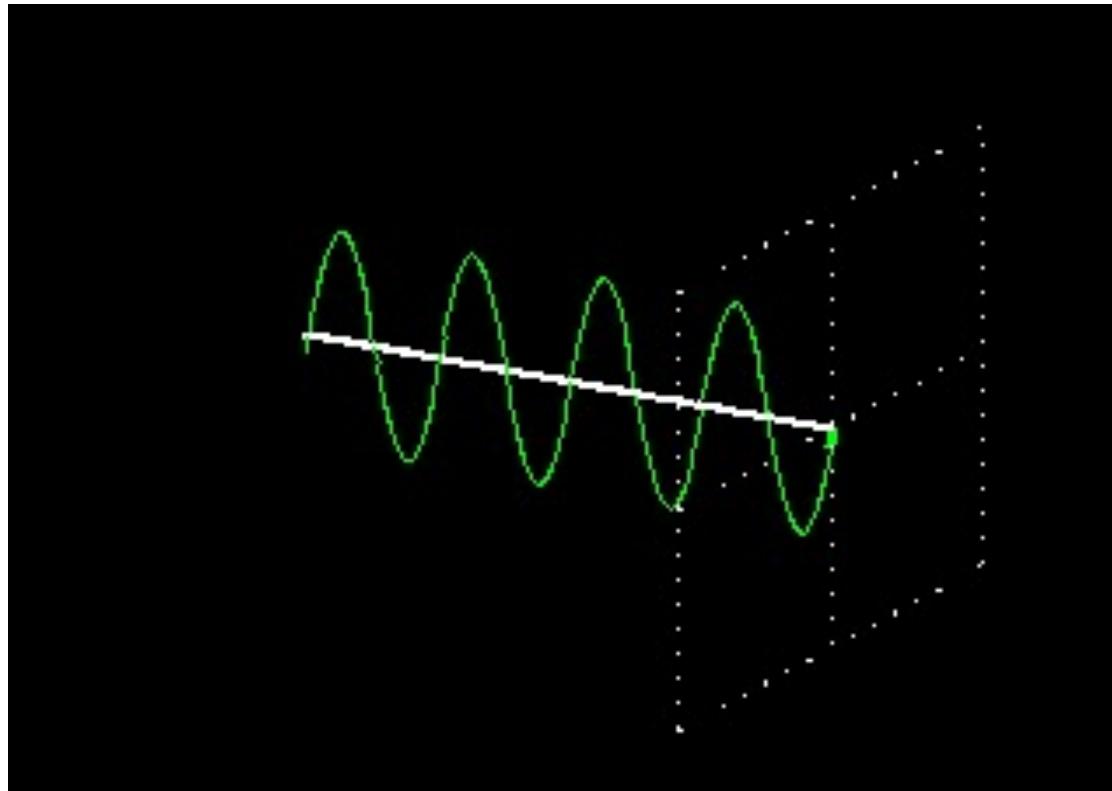
$$H(x,t) = H_{z0} \sin(2\pi ft + kx + \phi)$$

Wave vector: $|\vec{k}| = 2\pi / \lambda$ (1/m) **Wave number:** $1/\lambda$, cm^{-1}

*Most of the observable effects are due to
electric vector of EMW:*

$$I = E_{yo}^2, \text{ (W/m}^2\text{)}$$

Properties of Electric Waves



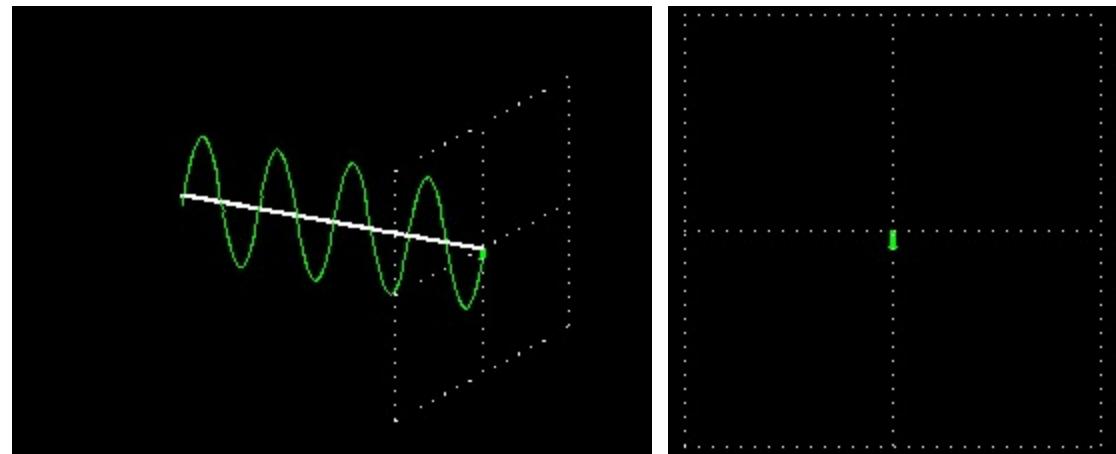
$$E(x,t) = E_{y0} \sin(2\pi ft + kx + \phi)$$

Linear/plane polarization

Plane-polarized electric wave

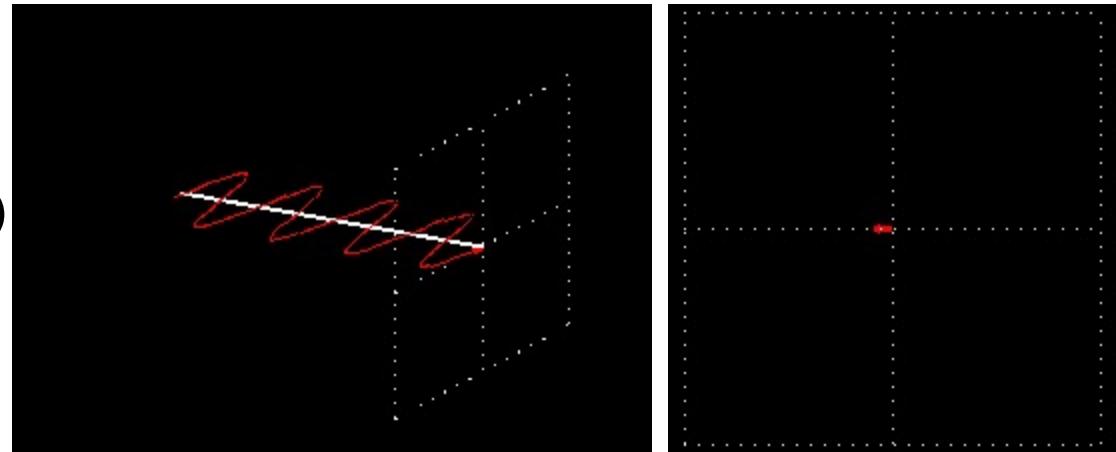
Vertically polarized

$$E_y(x,t) = E_{y0} \sin(2\pi ft + kx)$$

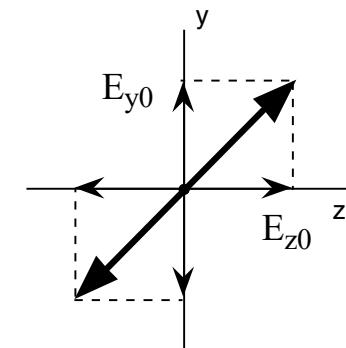
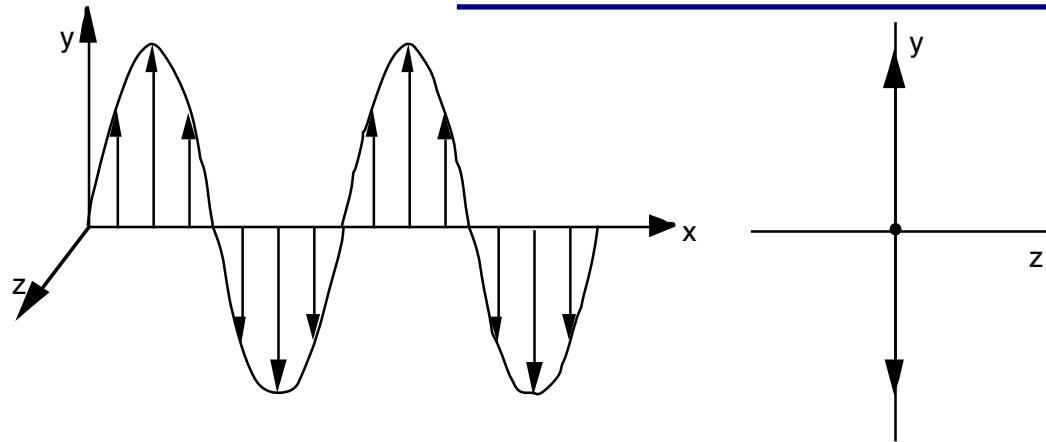


Horizontally polarized

$$E_z(x,t) = E_{z0} \sin(2\pi ft + kx)$$



POLARIZATION



$$\vec{E}(t, x) = \vec{E}_{y0} \sin(2\pi ft + kx + \phi_y) + \vec{E}_{z0} \sin(2\pi ft + kx + \phi_z)$$

What, if $E_{y0} = E_{z0}$ and $\delta = (\phi_y - \phi_z) = \pi/2 \pm n\pi$?

$$E^2 = E_0^2 \sin^2(2\pi ft + kx) + E_0^2 \sin^2(2\pi ft + kx + \pi/2)$$

$$E^2 = E_0^2 [\sin^2(2\pi ft + kx) + \cos^2(2\pi ft + kx)] = E_0^2$$

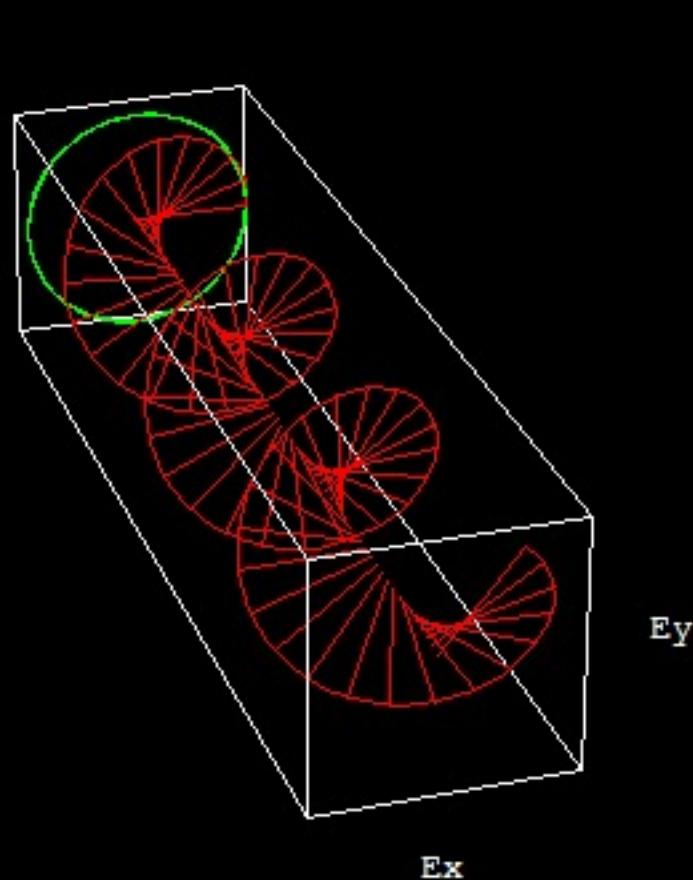
$$E_0^2 = x^2 + y^2 \leqslant \text{This is circle of radius } E_0: \text{Circular polarization}$$

Amplitude of electric vector is fixed

Tip of electric vector inscribes a circle

$$\delta = \pi/2 \pm n\pi$$

Circular polarization

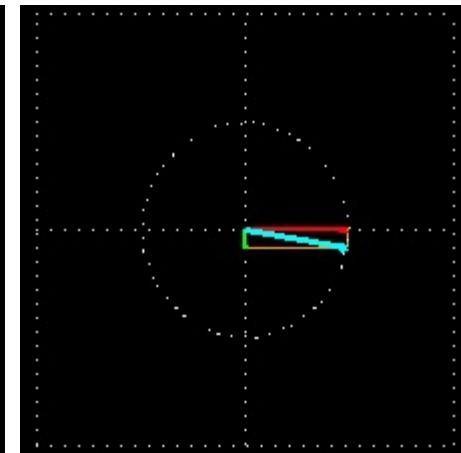
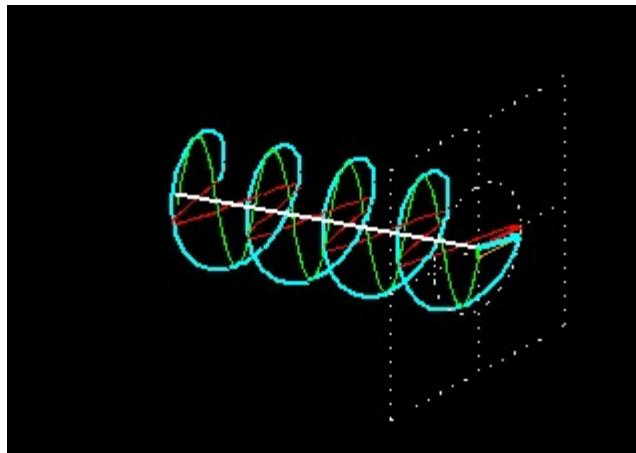


Circular polarized electric wave

Right circular

$$E_y(x,t) = E_{y0} \sin(2\pi ft + kx + 90^\circ)$$

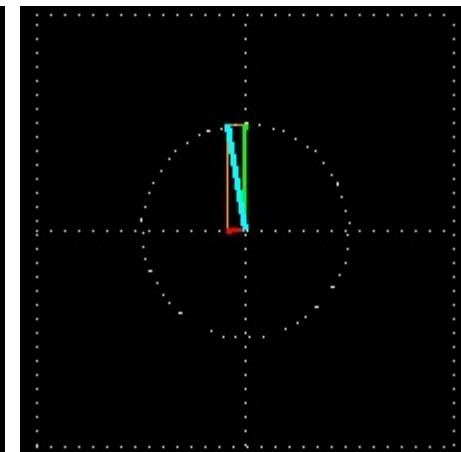
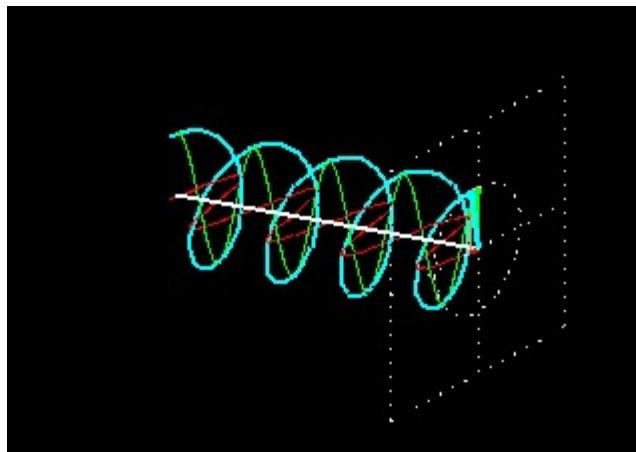
$$E_z(x,t) = E_{z0} \sin(2\pi ft + kx)$$



Left circular

$$E_y(x,t) = E_{y0} \sin(2\pi ft + kx - 90^\circ)$$

$$E_z(x,t) = E_{z0} \sin(2\pi ft + kx)$$



Light dualism

Light is just a particular case of EMW with the wavelengths, which are much smaller than the size of the object our vision can resolve.

According to the Photo-Electric Effect: light behaves like particles

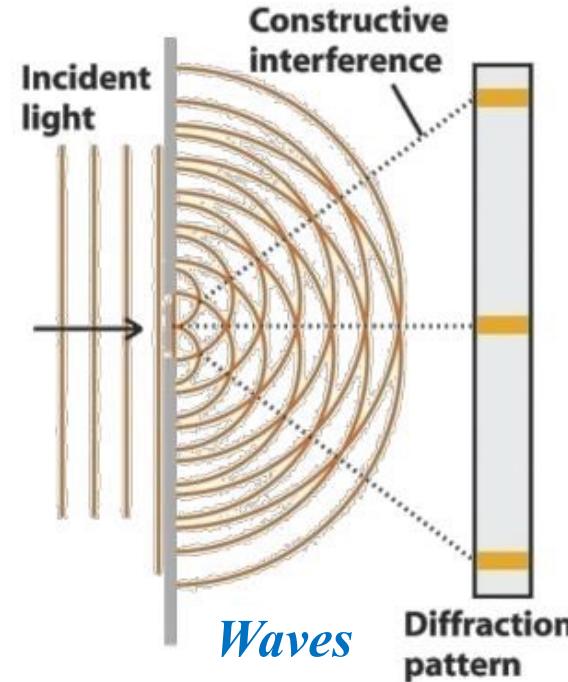
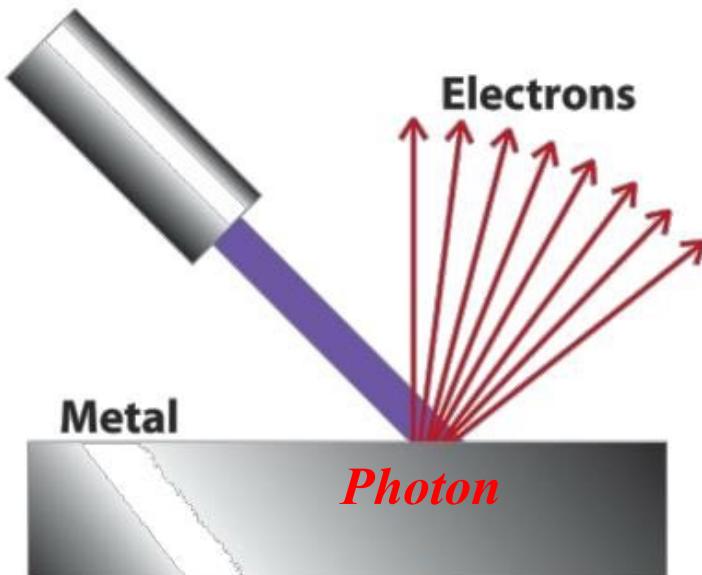
According to the 2-Slit Experiment: light behaves like waves

Interference of two waves

Modern concept of light:

Light has both wave-like and particle-like behavior

**Ultraviolet
radiation source**



Properties of photon

Photon energy: $E = h\nu$; $h = 6.626 \times 10^{-34}$ Joule·sec

Photon angular momentum (spin):

$J = \hbar/2\pi$; or in units of \hbar : $J=1 \Rightarrow$

photons are **Bosons**.

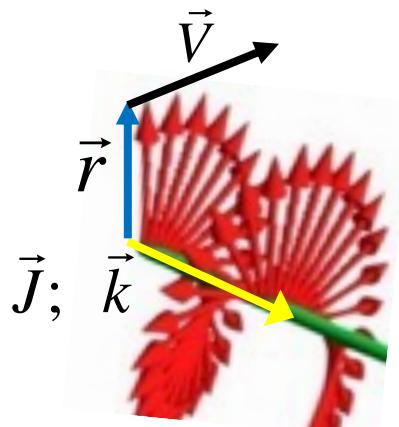
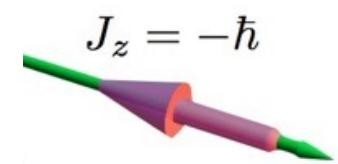
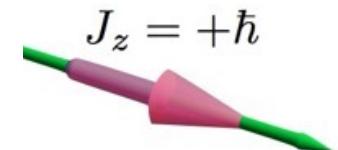
\vec{J} is a vector along the light propagation axis z :

$J_z = \pm 1$ corresponds to the **L**- and **R**- polarized light.

Linear polarized light is a linear combination of **L** and **R**.

$J_z \neq 0$, because photon is a **massless** particle that always moves with c .

Classical picture: $\vec{E} \perp \vec{k}$, but $\vec{J} = \vec{V} \times \vec{r} \Rightarrow \vec{J} \parallel \vec{k}$



Properties of photon

Photon energy: $E = h\nu$; $h = 6.626 \cdot 10^{-34}$ Joule·sec

Photon angular momentum (spin):

$J = \hbar = h/2\pi$; (h is Plank's constant), or in units of \hbar : $J=1 \Rightarrow$

photons are **Bosons**. \vec{J} is a vector along the light propagation axis Z :

$J_z = \pm 1$ corresponds to the **L**- and **R**-polarized light.

Linear polarized light is a linear combination of **L** and **R**.

Spectroscopic units:

Frequency ν : sec^{-1} or Hz

$$E = h\nu;$$

Wavelength λ : $\mu\text{m} = 10^{-6}\text{m}$ or

$$E = \frac{hc}{n \cdot \lambda_n}$$

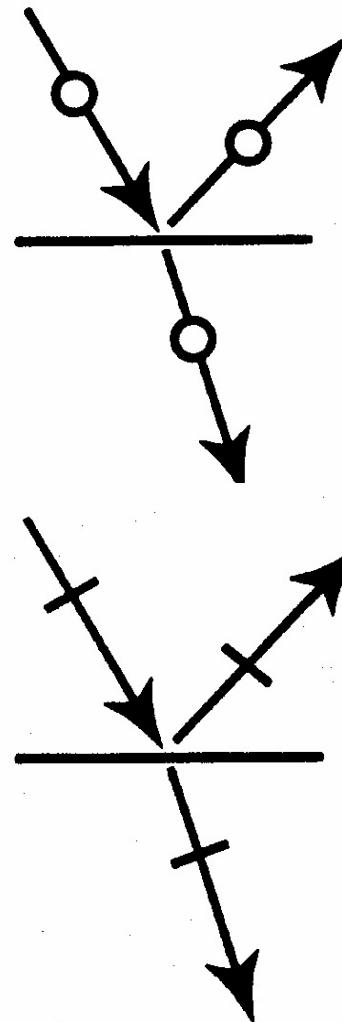
$\text{nm} = 10^{-9}\text{ m}$ or $\text{\AA} = 10^{-10}\text{m}$

Wavenumber: $\tilde{\nu} = \frac{1}{n \cdot \lambda_n} = \frac{\nu}{c}$ cm^{-1} only! $E = \frac{hc}{n\lambda_n} = hc\tilde{\nu}$

Reflection of polarized light

Perpendicular (“S”)
polarization **sticks** up out
of the plane of incidence.

Parallel (“P”) polarization
lies **parallel** to the plane of
incidence.



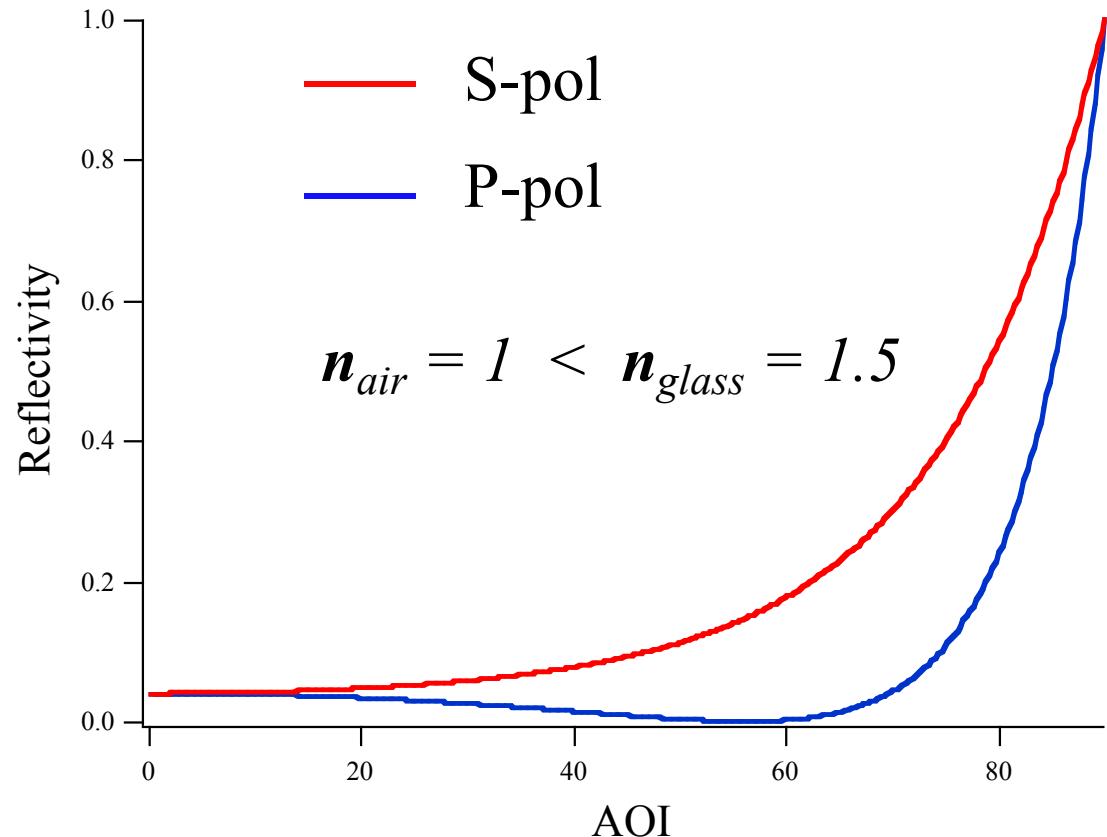
Does it matter for reflectivity ?

Reflection Coefficients for an Air-to-Glass Interface

Zero reflection for P polarization at
“Brewster's angle”
(56.3° here).

$$\tan \theta_{brewster} = \frac{n_t}{n_i}$$

$$R(0) = \frac{(n - 1)^2}{(n + 1)^2}$$

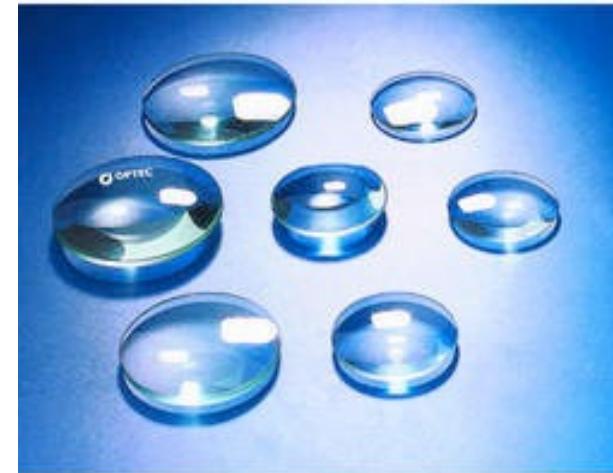


Reflectivity 4% at normal incidence

Making image with thin lens

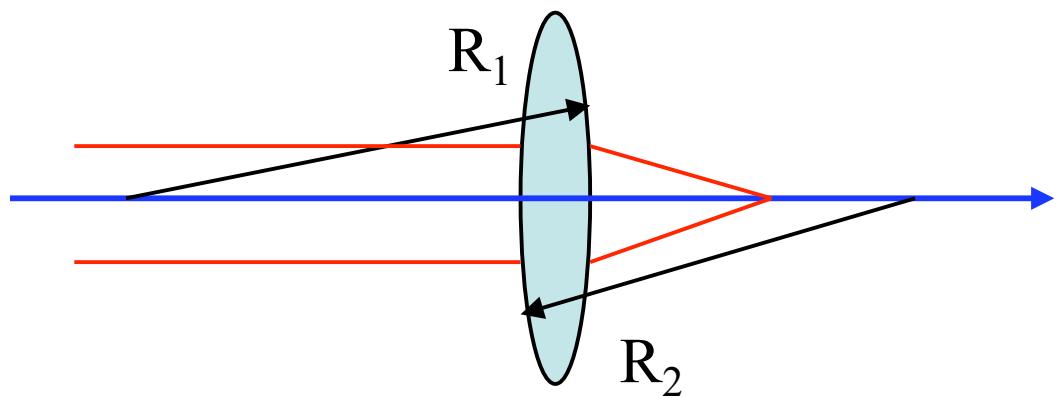
Lens is made of transparent material with n different from n of the media they are used in.

Glass, fused silica, plastic, ZnSe, Ge



They are optically polished (min scattering) with one or both sides curved (R_1, R_2); characterized by focal length f :

$$\frac{1}{f} = (n - 1) \left(\frac{1}{R_1} + \frac{1}{R_2} \right)$$



lens in air

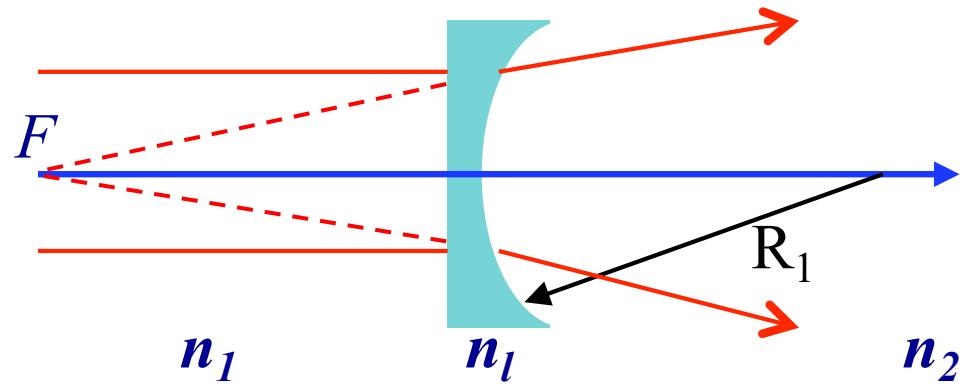
BI-convex lens: $R_1 > 0, R_2 > 0, \Rightarrow f > 0$; **positive**

Making image with thin lens

They are optically polished (min scattering) with one or two sides curved (R_1, R_2); characterized by focal length f :

$$\frac{1}{f} = (n - 1) \left(\frac{1}{R_1} + \frac{1}{R_2} \right)$$

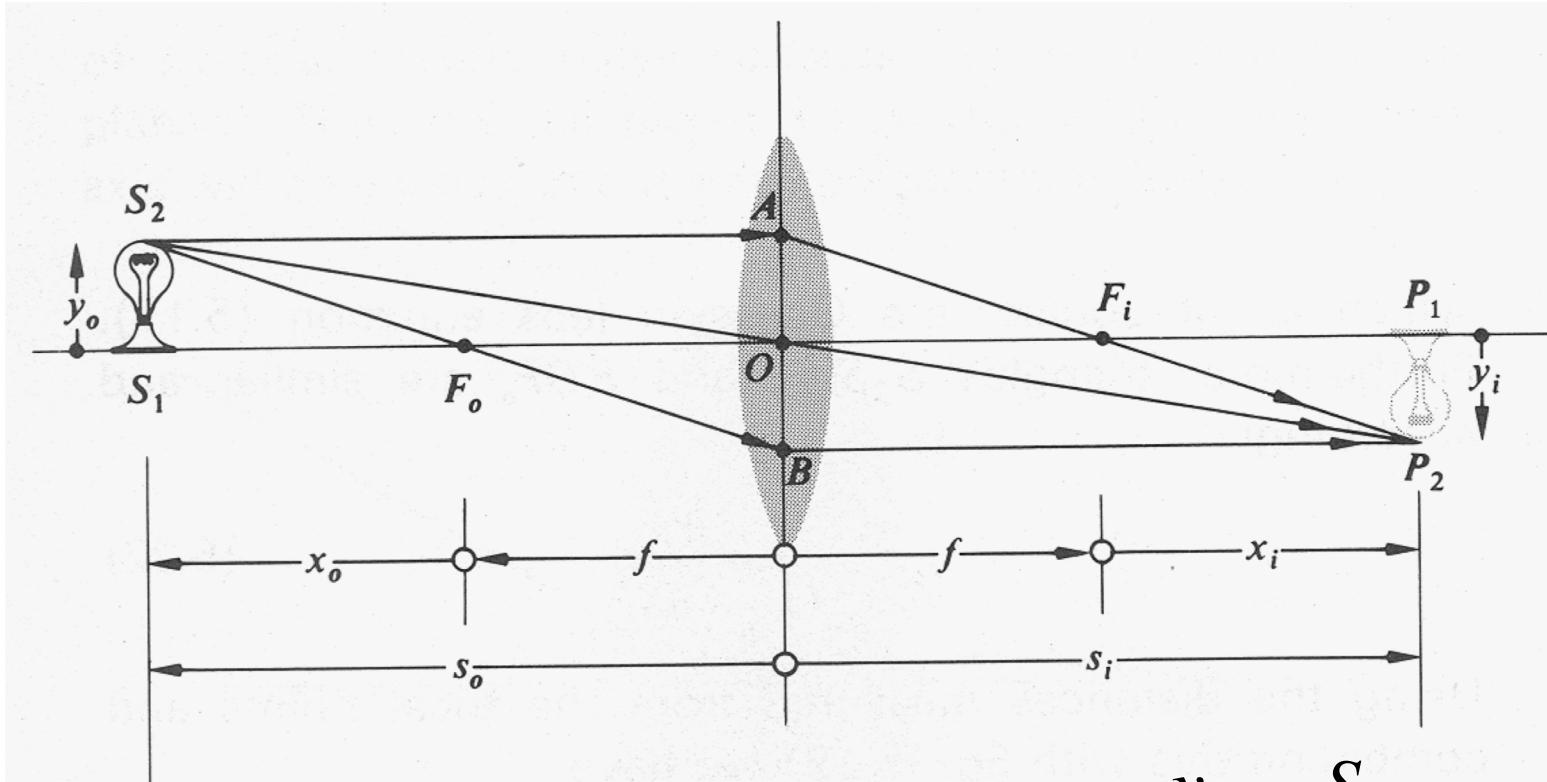
lens in air



Plano-concave lens: $R_1 < 0, R_2 = \infty \Rightarrow f < 0$; negative

Lens imaging:

$$\frac{1}{f} = \frac{1}{s_o} - \frac{1}{s_i} \quad \begin{array}{l} f \leftrightarrow \infty; \\ 2f \leftrightarrow 2f \end{array}$$



$$s_o > 0, \quad s_i < 0$$

$$M = \frac{y_i}{y_o} = \frac{s_i}{s_o}$$

<u>Name of Region</u>	<u>Wavelength Range</u>	<u>Wave number</u>	<u>Other Units</u>
	3000 m		10^5 Hz
Radio Waves	{ 0.3 m (30 cm)		10^9 Hz (1 GHz)
Microwaves (MW)	{ 0.06 cm (600 μ m)	16.6 cm^{-1}	
Far Infrared (FIR)	{ 30 μ m	333 cm^{-1}	
Mid Infrared	{ 3 μ m	3333 cm^{-1}	
Near Infrared (NIR)	{ 0.8 μ m (800 nm)	$12,500 \text{ cm}^{-1}$	1.55 eV
Visible (VIS)	{ 400 nm	$25,000 \text{ cm}^{-1}$	3.1 eV
Ultraviolet (UV)	{ 200 nm (2000 \AA)	$50,000 \text{ cm}^{-1}$	6.19 eV
Vacuum Ultraviolet (VUV)	{ 100 nm (1000 \AA)	$100,000 \text{ cm}^{-1}$	12.4 eV
Extreme Ultraviolet (XUV)	{ 5 nm (50 \AA)		247.8 eV
x-rays and γ rays	{ 10^{-3} \AA		

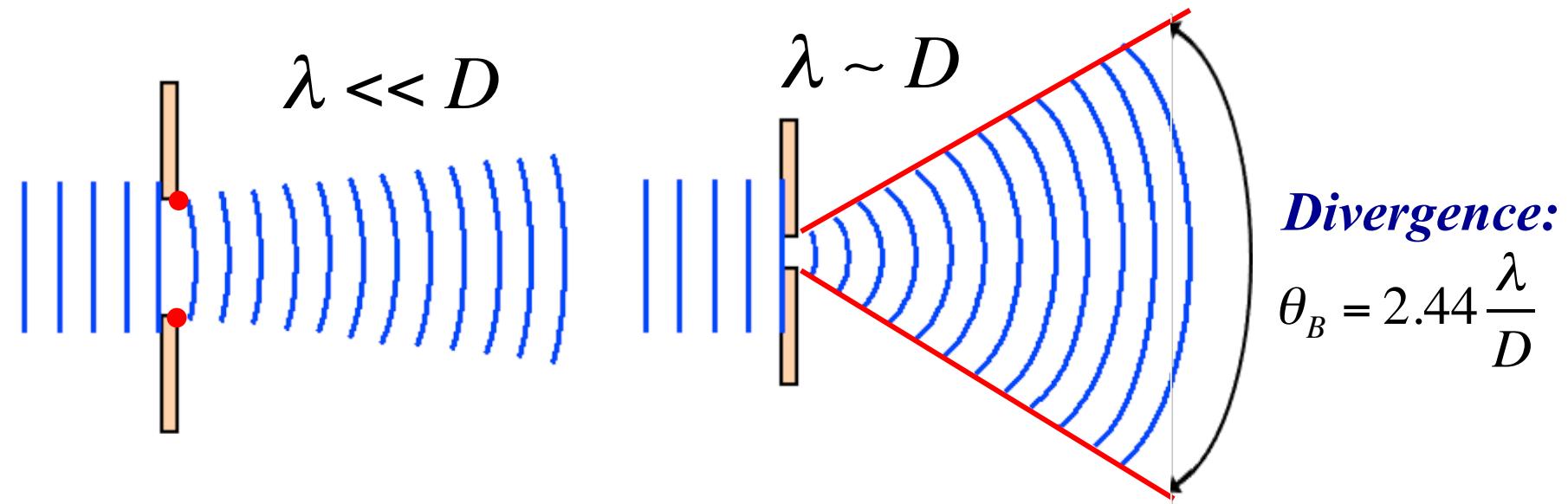
Wave divergence

Wave expands (diverges) upon propagation and becomes larger than the entry.

$$\lambda \sim D$$



Diffraction on hole for $\lambda \approx D$

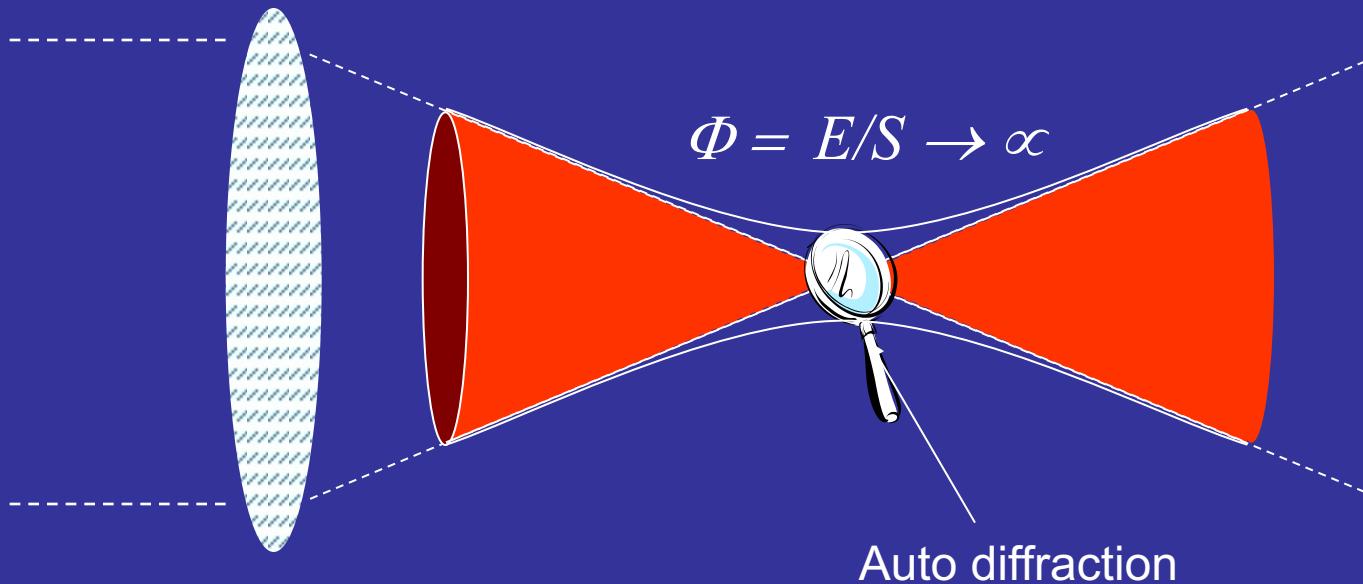


Any beam is divergent; at best, the divergence for a diffraction limited beam is $\theta = 2.44 \frac{\lambda}{D}$

Gaussian beams

Gaussian beams

What happens when we try to focus a laser beam?



The beam has a finite waist size !

Gaussian beams

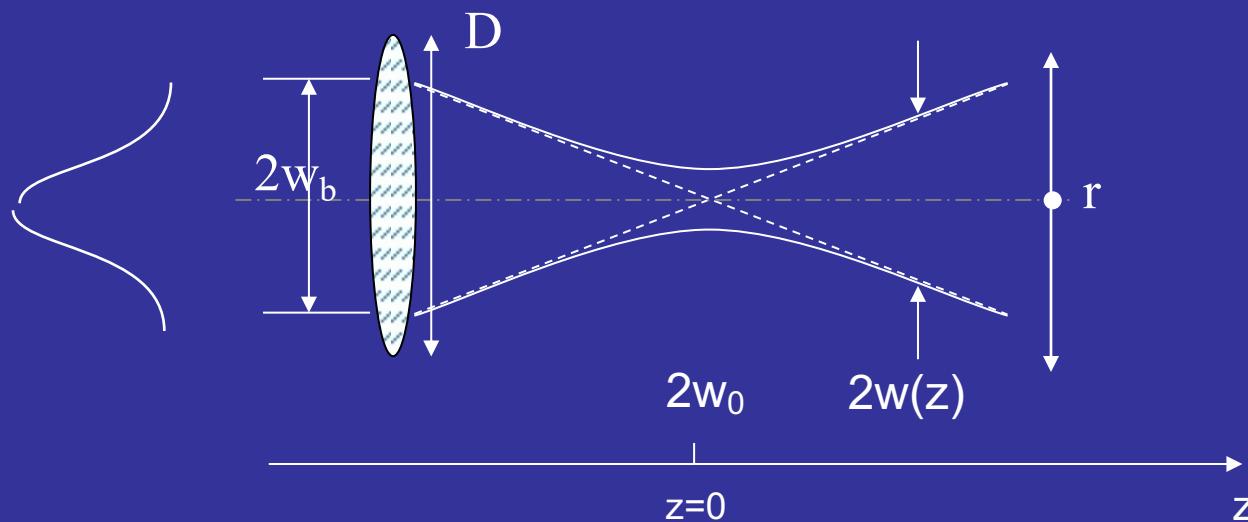
Hermite-Gaussian polynomials :

$$E_{mn}(x, y, z) = \frac{w_0}{w(z)} H_m\left(\frac{\sqrt{2}x}{w(z)}\right) H_n\left(\frac{\sqrt{2}y}{w(z)}\right) \exp\left[-\frac{r^2}{w^2(z)} - i\phi(r, z)\right]$$

Fundamental Gaussian beam ($m=0, n=0$):

Intensity:

$$I(r, z) = \frac{2P}{\pi w(z)^2} \exp\left(-2 \frac{r^2}{w(z)^2}\right)$$

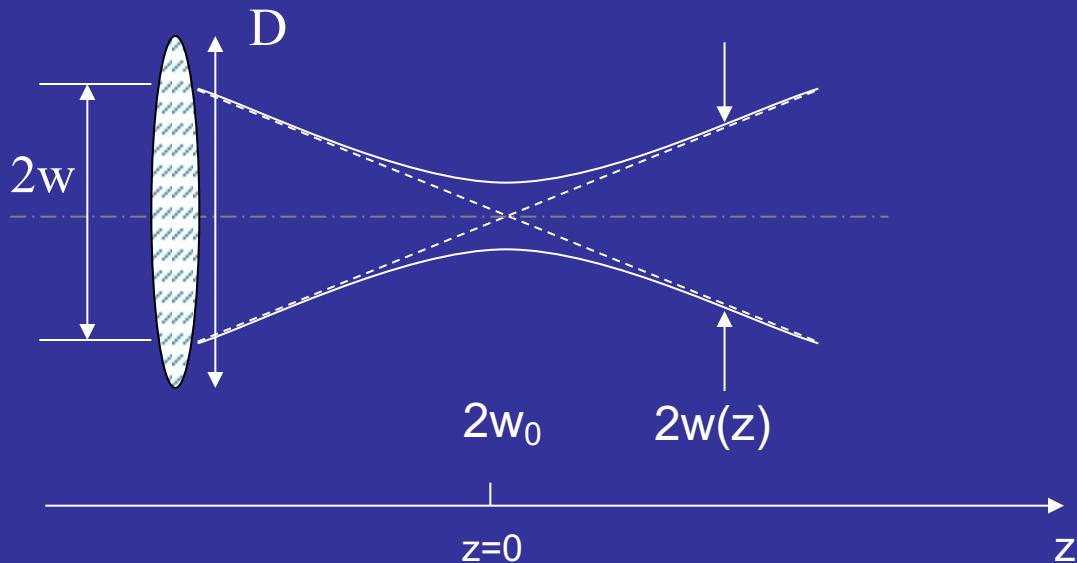


Fundamental Gaussian beam

Beam waist:

$$(2w < D)$$

$$2w_{\min} = 2.44 \frac{\lambda \cdot F}{2w}$$



Beam waist:

$$(2w \geq D)$$

$$2w_0 = 2.44 \frac{\lambda \cdot F}{D}$$

Divergence:

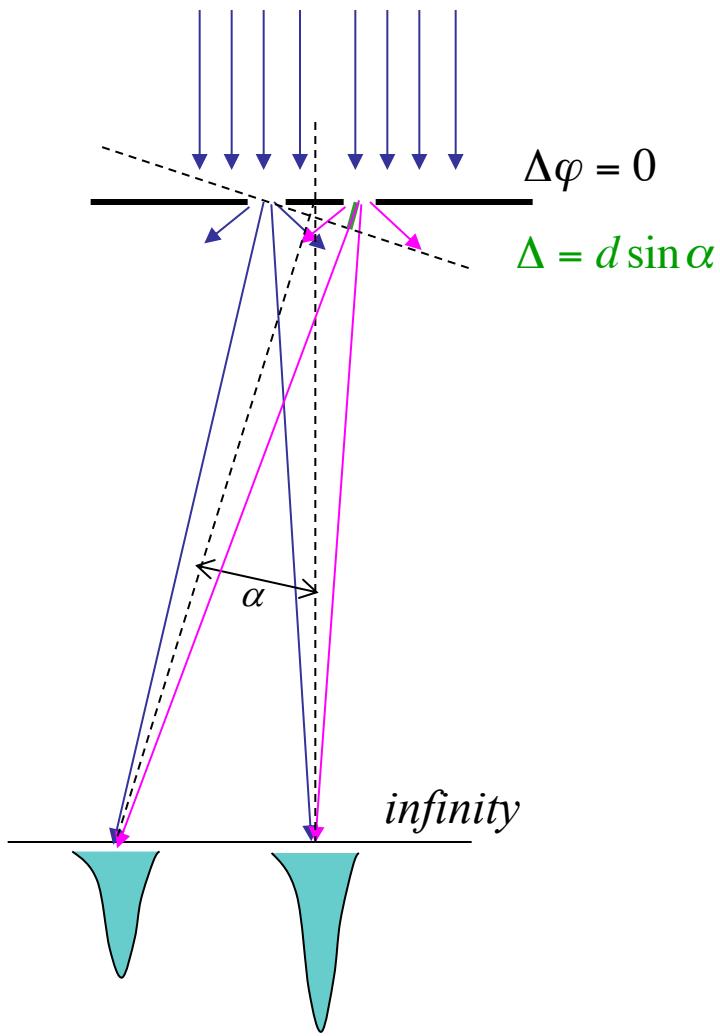
$$\theta_B = 2.44 \frac{\lambda}{D}$$

$$2\omega_0 = \theta_B \cdot F$$

Interferometric Optics

- *Diffraction grating*
- *Interferometer Fabri-Perout*

Interference

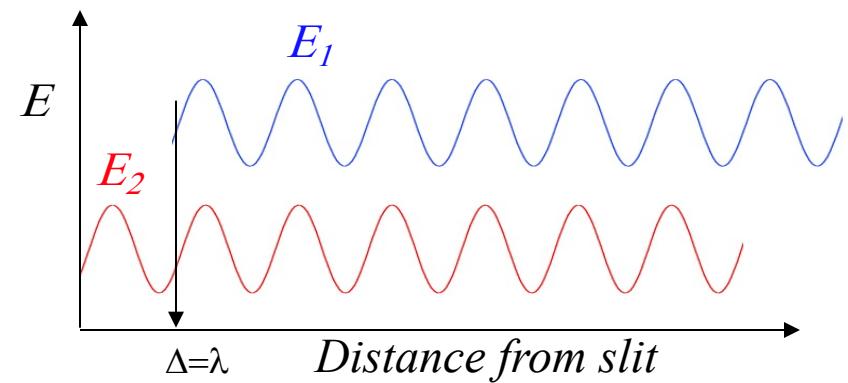


$$I = (\vec{E}_1 + \vec{E}_2)^2 = E_1^2 + E_2^2 + 2E_1E_2 \cos(\varphi)$$

For max: $\varphi_{\max} = 2\pi n; n=0, 1, 2, \dots$

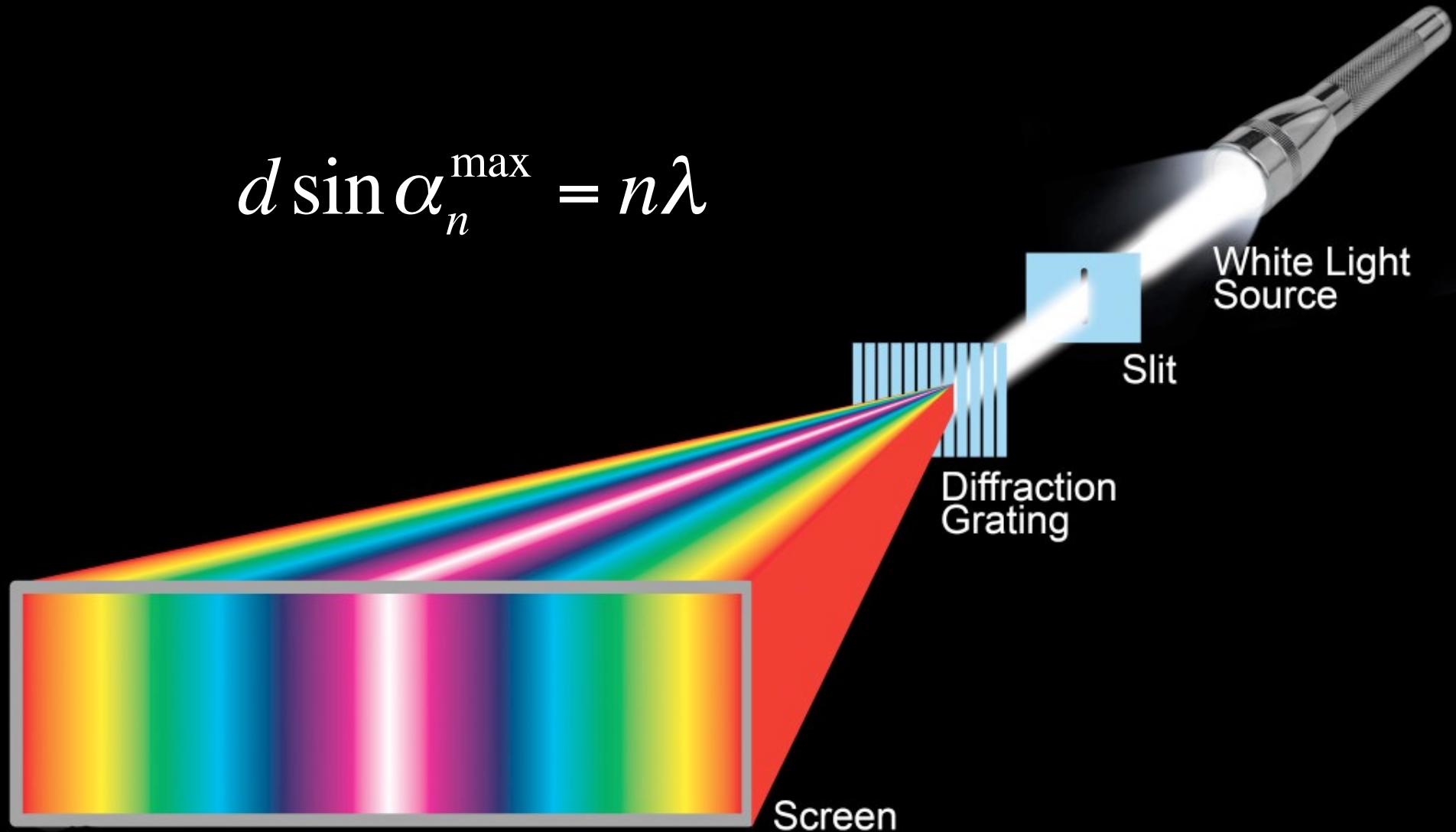
$$\varphi = 2\pi \frac{\Delta}{\lambda};$$

$$\Delta_{\max} = n\lambda$$



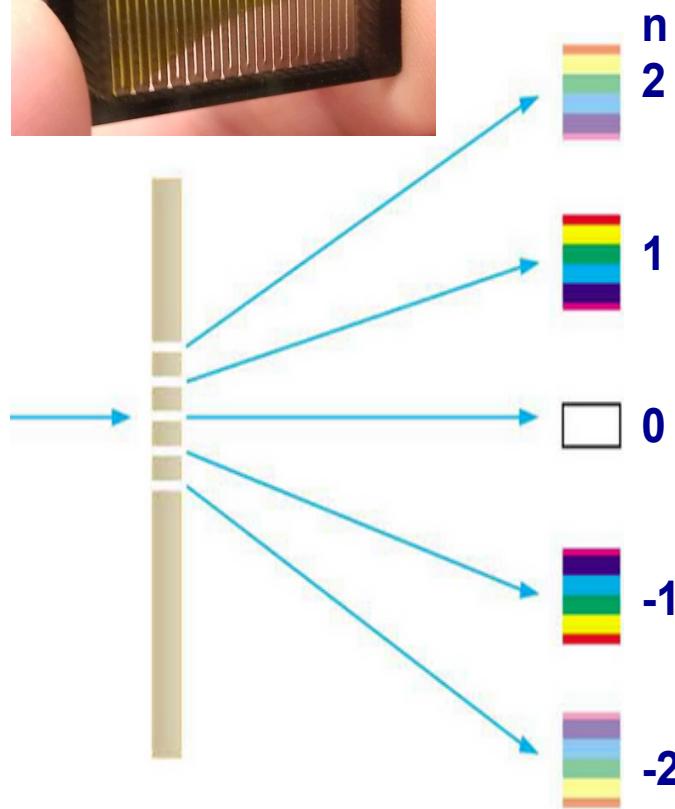
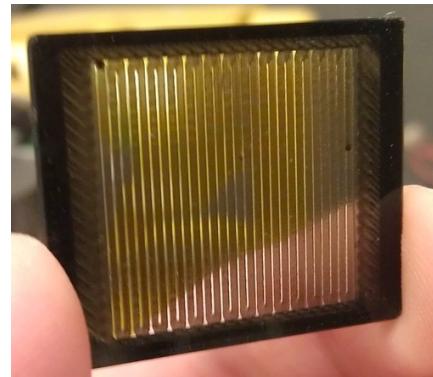
Multiple slit Interference

$$d \sin \alpha_n^{\max} = n\lambda$$



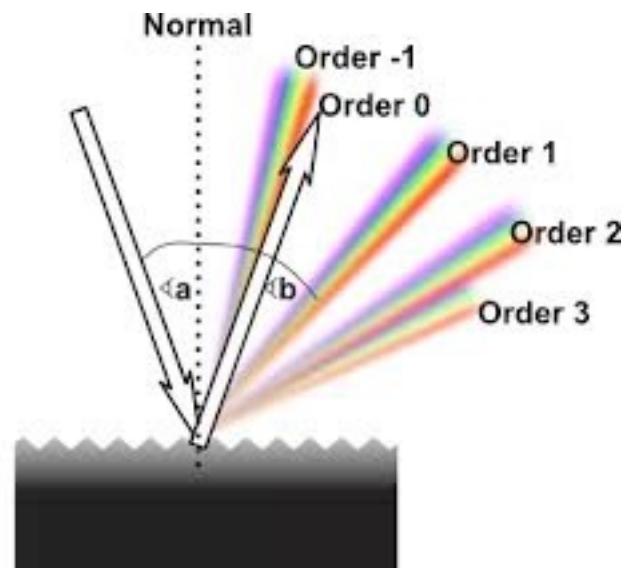
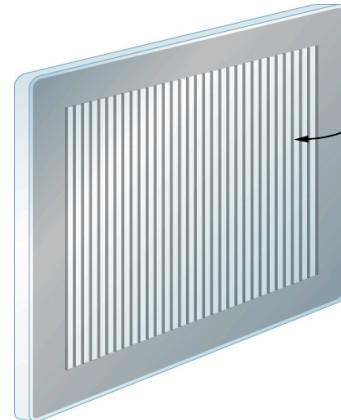
Diffracting grating

An optical device that disperses incoming rays of different wavelengths to different directions



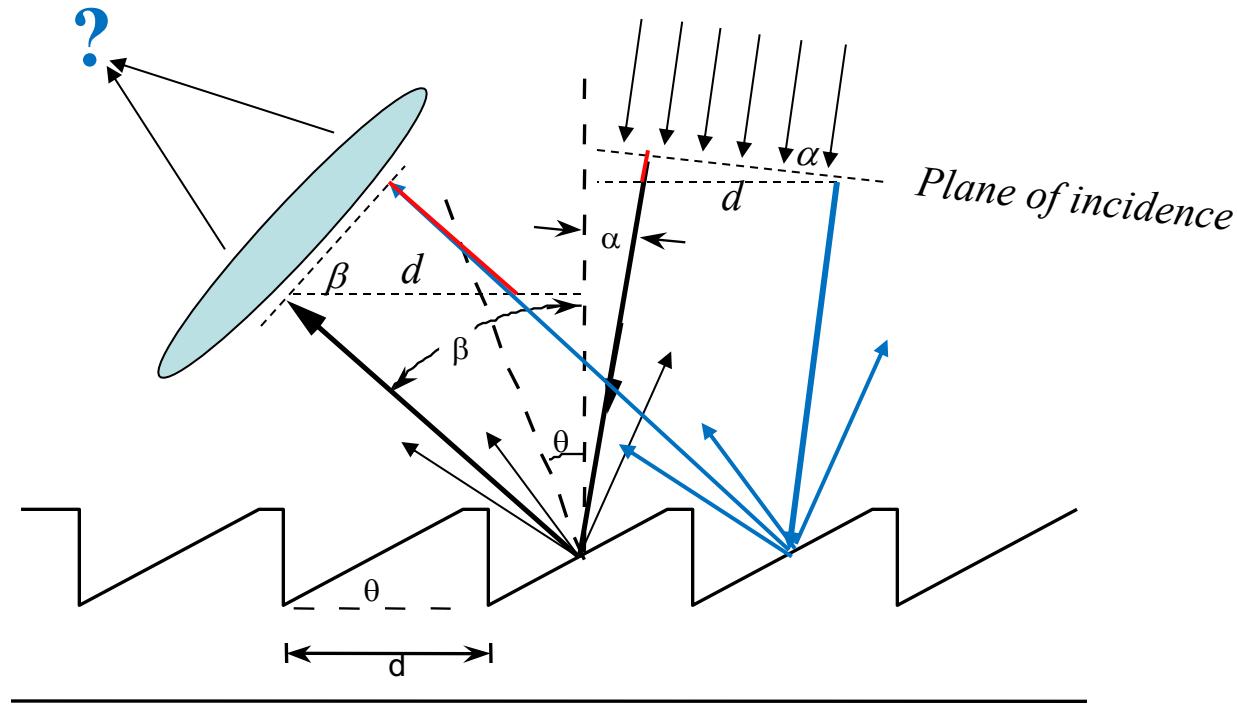
$$d \sin \alpha_n^{\max} = n\lambda;$$

$$\lambda < d$$



Interferometric Optics

Grating:



Path difference: $\Delta = d (\sin \alpha \pm \sin \beta)$

+ if α and β are on the same side;
 - If α and β are on opposite sides

Phase difference: $\varphi = 2\pi f \cdot \Delta t = 2\pi \frac{f \Delta}{c} = 2\pi \frac{\Delta}{\lambda}$

$$\varphi = 2\pi \frac{d(\sin \alpha \pm \sin \beta)}{\lambda}$$

Grating

Electrical field by each ray: $E_n = E_1 \cos(n\cdot\varphi)$

Total electrical field:

$$E_{tot} = \sum E_1 \{1 + \cos\varphi + \cos 2\cdot\varphi + \dots + \cos N\cdot\varphi\}$$

Total intensity: $I_{tot} = |E_{tot}|^2$

Complex function presentation

1. Euler formula:

$$e^{i\vartheta} = \cos\vartheta + i \sin\vartheta \Rightarrow \operatorname{Re}(e^{i\vartheta}) = \cos\vartheta$$

$$E_n = E_1(\cos n\varphi) = E_1 \operatorname{Re}(e^{in\varphi})$$

Grating

Total amplitude:

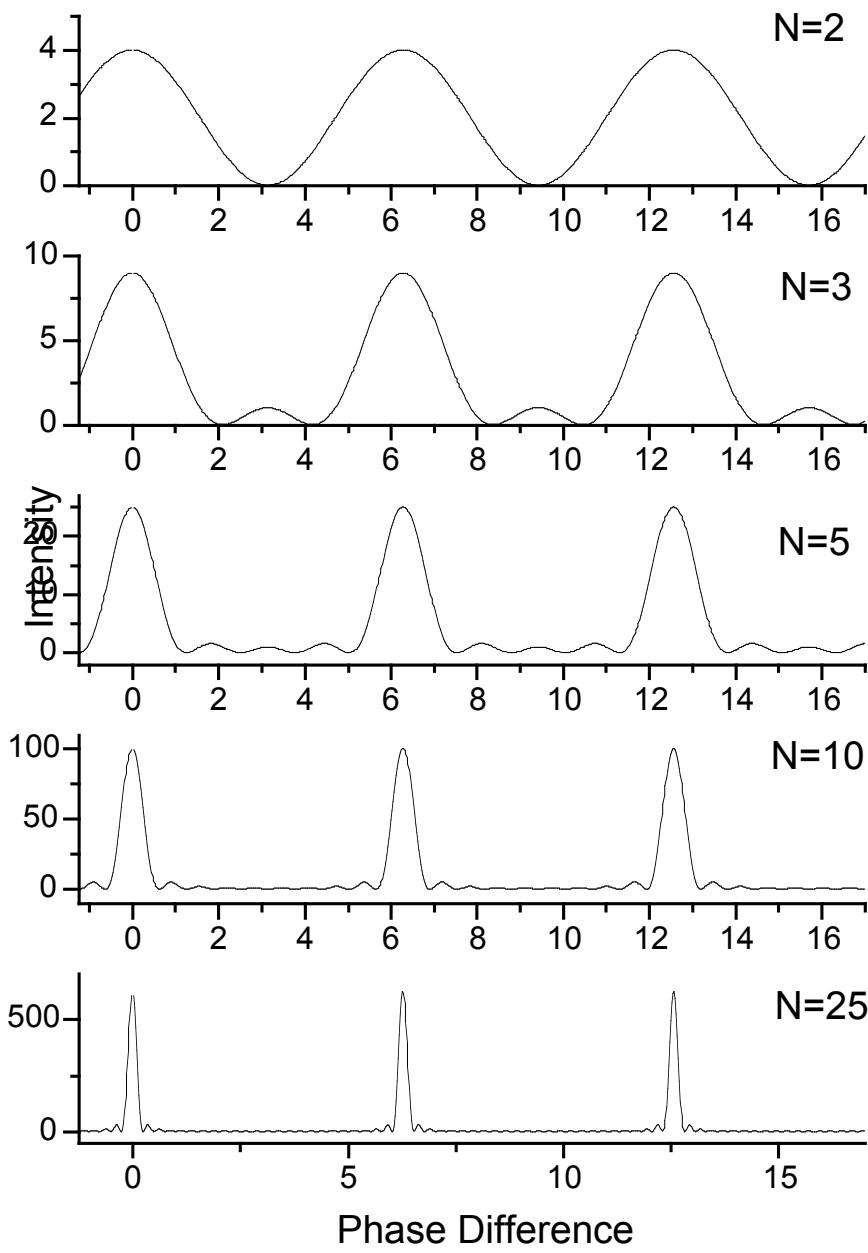
$$E_{tot} = E_1 (1 + e^{i\phi} + e^{i2\phi} + e^{i3\phi} + \dots + e^{iN\phi})$$

$$E_{tot} = E_1 (1 + e^{i\phi} + e^{i2\phi} + e^{i3\phi} + \dots + e^{iN\phi}) = E_1 \frac{1 - e^{iN\phi}}{1 - e^{i\phi}}$$

Total intensity:

$$I_{tot} = |E_{tot}|^2 = E_0^2 \frac{1 - \cos N\phi}{1 - \cos \phi} = I_0 \frac{\sin^2(N\phi/2)}{\sin^2(\phi/2)} \simeq I_0 N^2$$

Total intensity:



$$I_{tot} = I_0 \frac{\sin^2(N\varphi/2)}{\sin^2(\varphi/2)} \rightarrow I_0 \cdot N^2; \text{ for } (\varphi/2 \rightarrow 0 + n\pi);$$

Maximum intensity:

$$\varphi/2 = 0, \pi, 2\pi, 3\pi, \dots m\pi$$

$$\varphi = 2\pi \frac{d(\sin \alpha \pm \sin \beta)}{\lambda}$$

or: $\pi \frac{d(\sin \alpha \pm \sin \beta)}{\lambda} = m\pi \quad m = 0, \pm 1, \pm 2$

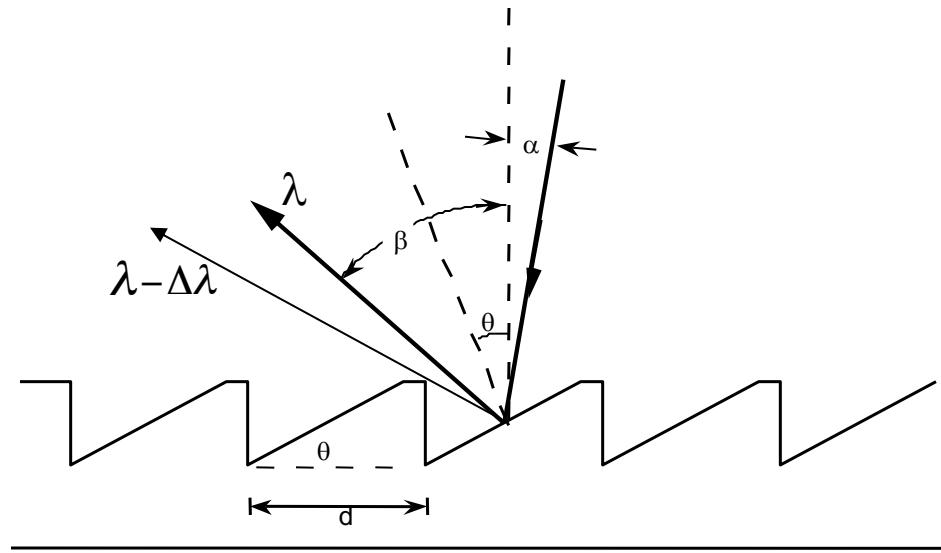
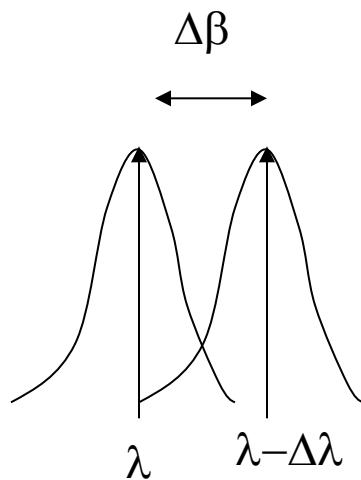
or: $m\lambda = d(\sin \alpha \pm \sin \beta) \quad m = 0, \pm 1, \pm 2$

$$\lambda \leq 2d$$

Minimum intensity:

$$(m + \frac{1}{N})\lambda = d(\sin \alpha \pm \sin \beta) \quad m = 0, \pm 1, \pm 2$$

Relay resolution criteria:



Max: $m\lambda = d(\sin \alpha \pm \sin \beta);$

Min: $(m + \frac{1}{N})(\lambda - \Delta\lambda) = d(\sin \alpha \pm \sin \beta);$

$$\cancel{m\lambda = m\lambda - m\Delta\lambda + \frac{\lambda}{N} - \frac{\Delta\lambda}{N}}$$

$$\frac{\lambda}{\Delta\lambda} = N \left(m + \frac{1}{N} \right)$$

$$m \gg \frac{1}{N},$$

$$\frac{\lambda}{\Delta\lambda} = mN = m\rho l$$

<i>Type</i>	<i>Capacity</i>	<i>Track pitch</i>	<i>Wavelength of laser light</i>
<i>CD</i>	<i>0.7 GB</i>	<i>1.6 μm</i>	<i>780 nm (0.780 μm)</i>
<i>DVD</i>	<i>4.7 GB</i>	<i>0.74 μm</i>	<i>650 nm (0.650 μm)</i>
<i>Blu-ray Disc</i>	<i>25 GB</i>	<i>0.32 μm</i>	<i>405 nm (0.405 μm)</i>

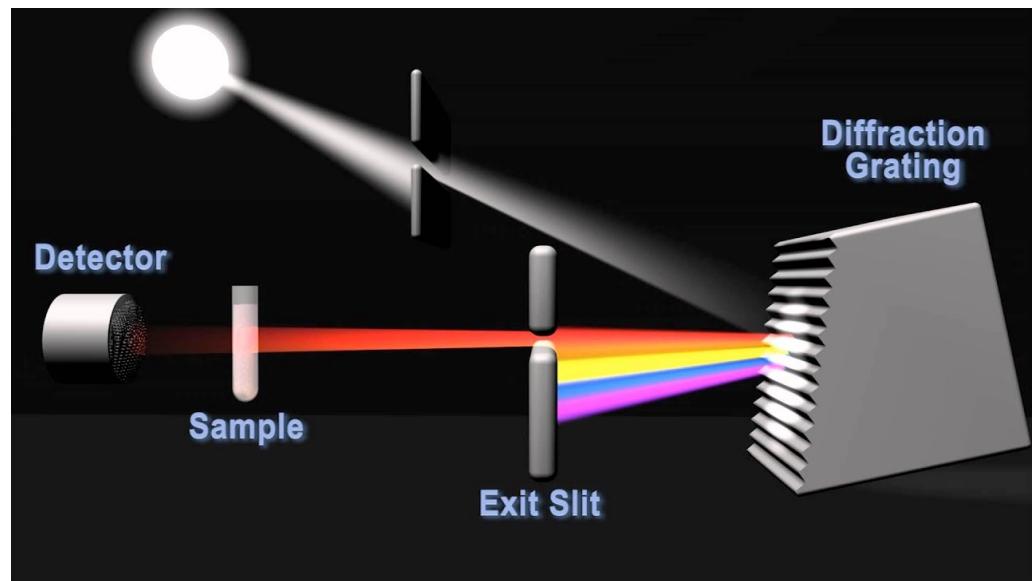
$$m\lambda = d(\sin \alpha \pm \sin \beta) \Rightarrow \lambda = \frac{2d}{m};$$

$$\lambda_1 < 2d, \lambda_2 < d \dots;$$

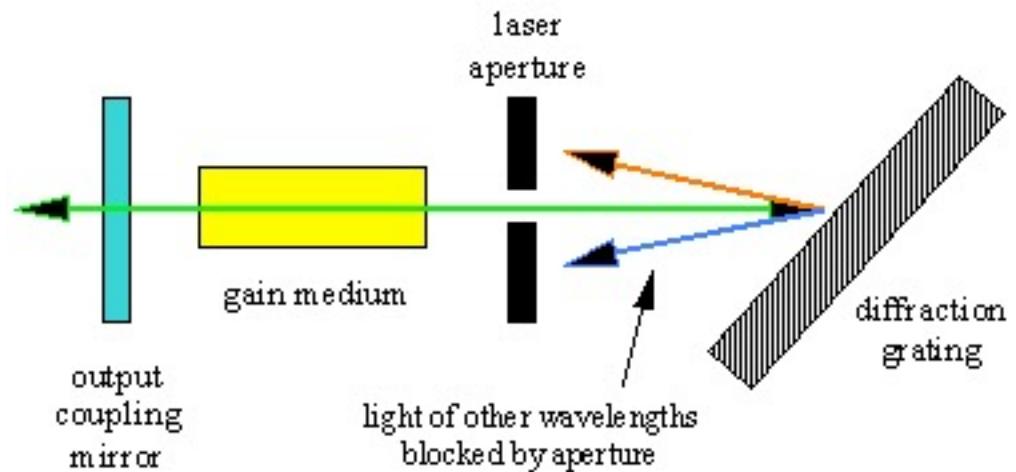
$$m_{\max} \leq \frac{2d}{\lambda}.$$

Diffraction grating

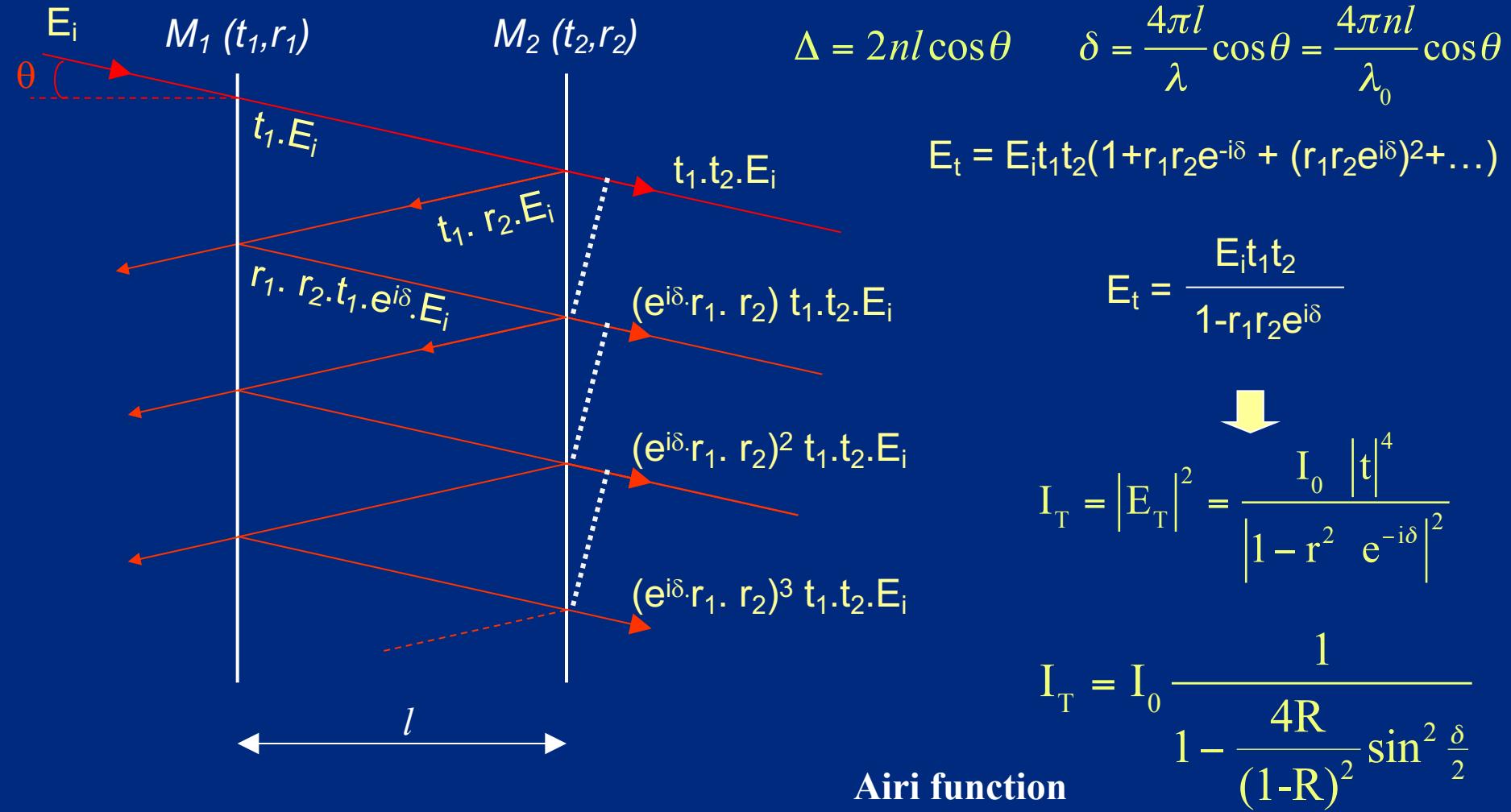
- *Analyze incoming light (spectrometer).*
- *Filter light of one particular wavelength (monochromator).*



- *Controlling light wavelength in tunable lasers (grating = selective mirror)*

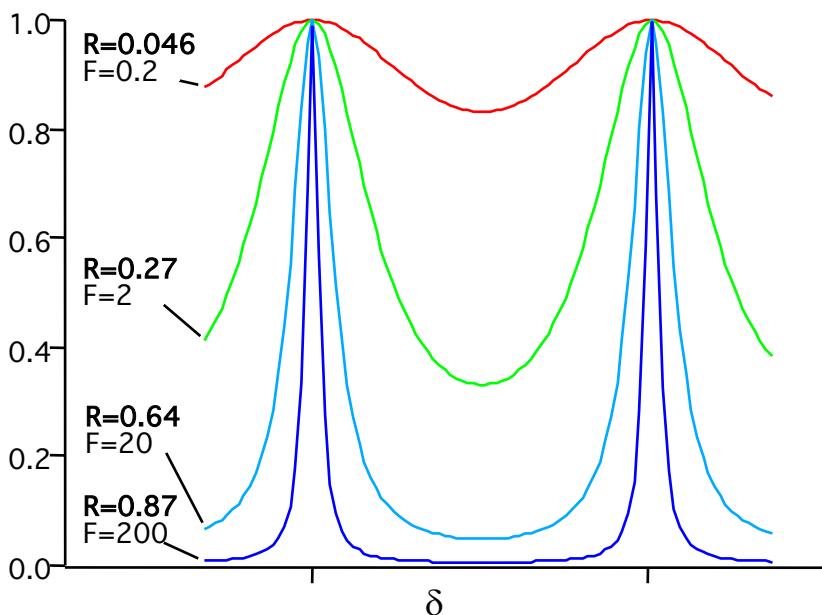


Fabry-Perot interferometer



Fabry-Perot interferometer

Airy function



$$I_T = I_0 \frac{1}{1 - \frac{4R}{(1-R)^2} \sin^2 \frac{\delta}{2}}$$

$$\delta = \frac{4\pi l}{\lambda} \cos \theta$$

Max: $\delta = 2m\pi, \rightarrow$ or: $2l \cos \theta = m\lambda$

Max transmittance: $\nu_m = \frac{mc}{2l \cos \theta}, \tilde{\nu}_m = \frac{m}{2l \cos \theta}$

$$\nu_m = c/\lambda_m$$

$$m = 0, 1, 2, \dots$$

Free spectral range: $\Delta \nu = \frac{c}{2l}; \Delta \tilde{\nu} = \frac{1}{2l};$

Transmission width
(and resolution): $\Delta \nu_{1/2} = \frac{\Delta \nu}{F} = \frac{(1-R)}{2\pi \cdot l \cdot \sqrt{R}}$

Finesse: $F = \frac{\pi \sqrt{R}}{(1-R)}$

FPI is a filter that transmits light with the frequencies multiple to the FSP:

$$\nu_m = m \cdot \frac{c}{2l} \cos \theta$$

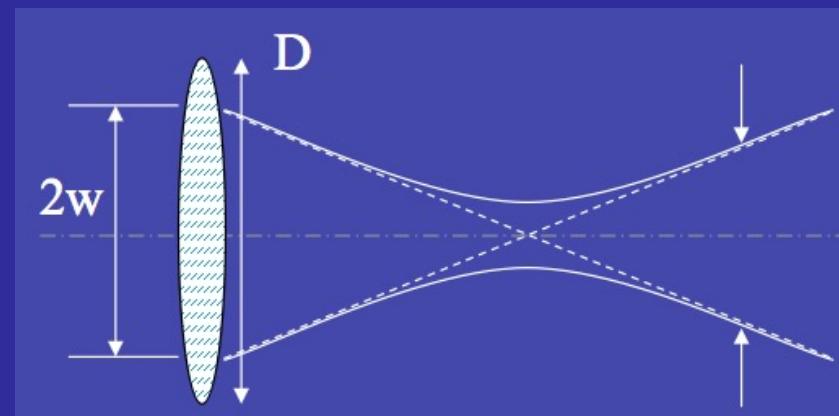


Summary for this lecture

Beam divergence: $\theta_B = 2.44 \frac{\lambda}{D}$

The smallest focal point:

$$2w_0 = 2.44 \frac{\lambda \cdot F}{D}$$



Intensity for N interfering rays of light:

$$I = I_0^2 \frac{\sin^2(N\varphi/2)}{\sin^2(\varphi/2)}$$

$$d \sin \alpha_n^{\max} = n\lambda$$

Spectral resolution of grating:

$$\frac{\lambda}{Nd} = \frac{n\Delta\lambda}{d} \Rightarrow \frac{\lambda}{\Delta\lambda} = Nn$$

Transmittance of IFP:

$$v_m = m \cdot \frac{c}{2l} \cos \theta$$